

in $[\text{Cu}(\text{bipy})_2(\text{ONO})]\text{NO}_3$. The e.s.r. data on $\text{Cu}(\text{en})_3\text{SO}_4$ are unusual in that magnetically non-equivalent sites can be distinguished (by three signals) in a non-dilute copper(II) complex.

A second case of a room-temperature dynamic Jahn–Teller effect is postulated^{5,6} in $\text{K}_2\text{Pb}[\text{Cu}(\text{NO}_2)_6]$ and here again “trapping-out” appears at lower temperatures. In view of the point made above, it is worth noting that separate signals are again found in a non-dilute copper(II) complex, again with no observable hyperfine structure (hyperfine structure has been noted in some non-dilute copper complexes, for example⁷ $\text{Cu}(\text{NO}_3)_4^{2-}$ and $\text{Cu}(\text{NCO})_4^{2-}$). Paoletti *et al.* have given a qualitative interpretation of partial data on a single crystal. They suggest the existence of two magnetically non-equivalent sites, implying that the unit cell of $\text{K}_2\text{Pb}[\text{Cu}(\text{NO}_2)_6]$ contains three molecules, since the intensity of site A is twice that of site B. However, the authors venture the unlikely suggestion that at low temperature one site (B) is distorted while the others remain regular octahedral.

A more plausible explanation is that *all* sites are distorted in the same way, their orientations conspiring to produce the observed effects in the rotation of the crystal. This would be feasible if two of the sites were inclined at identical angles to the rotation axis. Suppose the principal molecular axes (z) of both sites A are inclined at angle θ to the rotation axis (R) and are coplanar with this axis, as in Figure 2. Then the signals due to the 2A sites would vary in the same way, from g_{\perp} to $(g_{\parallel}^2 \sin^2 \theta + g_{\perp}^2 \cos^2 \theta)^{1/2}$; if θ is small then $g(\text{A})$ varies little, as observed. Now if the z axis of site B is orientated perpendicularly to the plane containing the $z(\text{A})$ axes, the observed spectra can be accounted for in that the B signal varies over the larger range of g_{\perp} to g_{\parallel} . In particular, the maximum value is shown for site B when sites A exhibit the minimum value (and *vice versa*), as observed. The attraction of this scheme is that it is based on the “trapping-out” of sites orientated in three mutually perpendicular directions, as typified by the $\text{Cu}(\text{en})_3\text{SO}_4$ case; the difference here is that $2\theta < 90^\circ$ in accordance with the lower crystal symmetry at this temperature.⁸

Although Paoletti *et al.* give no g -value data, the apparent absence of exchange-coupling (witness the

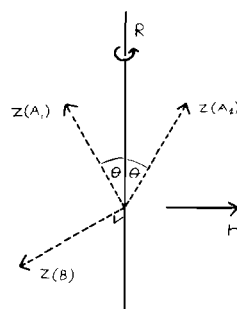


Fig. 2. A possible orientation, relative to the crystal rotation axis, of the principal axes of the three sites present in the unit cell of $\text{K}_2\text{Pb}[\text{Cu}(\text{NO}_2)_6]$ at low temperatures.

separate signals from the sites observed in the single-crystal data) suggests that the powder data of Elliott *et al.*⁸ may correspond with the molecular g -values. We then have values at 77 K of $g_{\parallel} = 2.061$ and $g_{\perp} = 2.155$, which suggest again the presence of a predominantly d_{z^2} ground state. However, the trapped distortion cannot have an exact axial symmetry, otherwise g_{\parallel} would equal 2.002. We suppose then that the symmetry is derived from a rhombic perturbation of a basically compressed tetragonal geometry. It then seems likely that the electronic reflectance spectrum³ can be assigned with the broad weak band at 7.0 kK ascribed to the $d_{x^2 - y^2} \rightarrow d_{z^2}$ transition with the other three transitions superposed under the sharper peak at 16.5 kK.

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